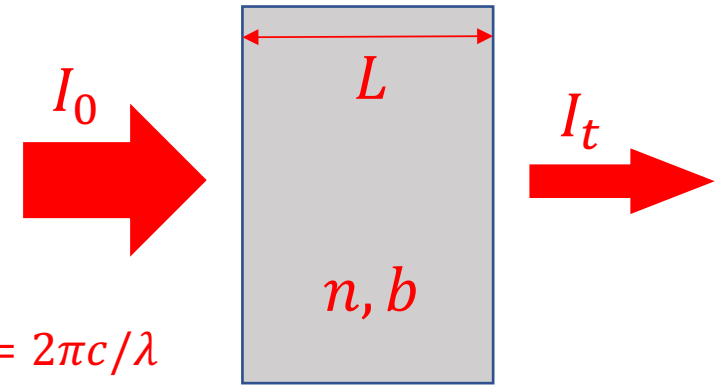


# Forward Scattering of Light: Cooperativity, Superflash, and More

**David Wilkowski**

Beer-Lambert's law:  $I_t(\omega) = I_0(\omega)e^{-b(\omega)}$



Optical frequency:  $\omega = kc = 2\pi c/\lambda$

For the electric field:

$$E_t(\omega) = E_0(\omega)\exp[in(\omega)kL] = E_0(\omega)\exp\left[-\frac{b(\omega)}{2} + i\phi(\omega)\right]\exp(ikL)$$

With a two-level atom (RWA):  $n(\omega) = 1 + \rho\alpha(\omega)/2$

In a dilute medium  $\rho\lambda^3 \ll 1$

Atomic polarizability

$$\alpha(\omega) = -\frac{3\pi}{k^3} \frac{\Gamma}{\delta + i\Gamma/2}$$

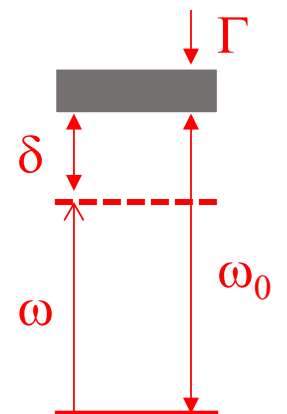
$$b = \text{Im}\{\alpha\}\rho kL$$

$$\phi = \frac{\text{Re}\{\alpha\}}{2}\rho kL$$

Optical  
theorem

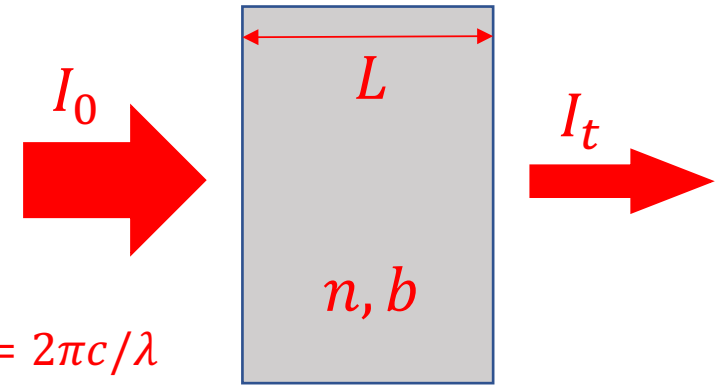
Scattering  
cross section

$$\delta = \omega - \omega_0$$



Beer-Lambert's law:  $I_t(\omega) = I_0(\omega)e^{-b(\omega)L}$

Optical frequency:  $\omega = kc = 2\pi c/\lambda$



For the electric field:

$$E_t(\omega) = E_0(\omega)\exp[in(\omega)kL] = E_0(\omega)\exp\left[-\frac{b(\omega)}{2} + i\phi(\omega)\right]\exp(ikL)$$

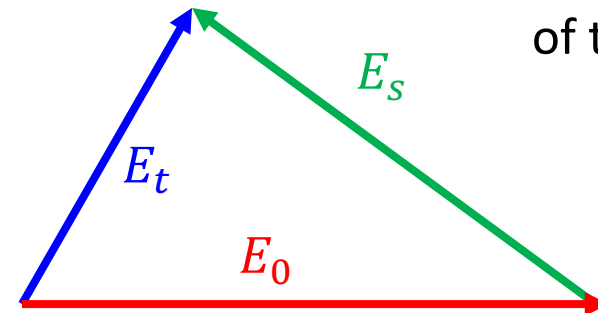
Coherent superposition of scalar fields:

$$E_t(t) = E_0(t) + E_s(t)$$

$E_t$ : the transmitted field

$E_0$ : the incident field

$E_s$ : the forward scattering field



Geometrical representation  
of the fields in the complex plane



Direct Observation of the Forward Scattering Field: Flash Effect

Superflash Effect:  $I_s|_{\text{sta}} > I_0|_{\text{sta}}$

Cooperative Pulse Train

Flash in  $\Lambda$ -Scheme

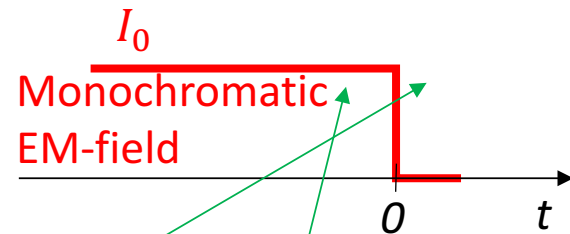
Outlook



Magneto-optical trap:  $10^9$  strontium atoms

# Direct Observation of the Forward Scattering Field

$$E_t(t) = E_0(t) + E_s(t) \rightarrow E_t(t) = \cancel{E_0(t)} + E_s(t)$$



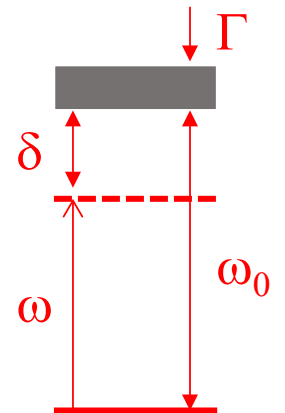
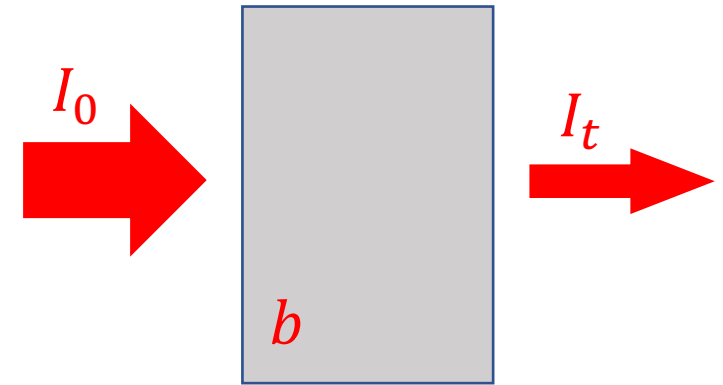
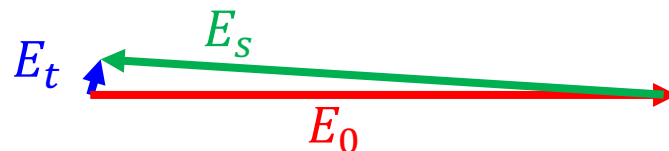
$$E_s(0^+) = E_s \Big|_{\text{sta}} \equiv E_s = E_t - E_0 = E_0 \left[ \exp \left( -\frac{b}{2} + i\phi \right) - 1 \right]$$

Medium response time has to be not “too fast”!

Atoms with large resonance quality factor:  $\frac{\Gamma}{\omega_0} \sim 10^{-11}$

If  $\delta \gg \Gamma \rightarrow (\phi, b) \sim 0 \rightarrow E_s \sim 0$

If  $\delta \ll \Gamma$  and  $b \gg 1 \rightarrow E_t = 0, E_s \sim -E_0$



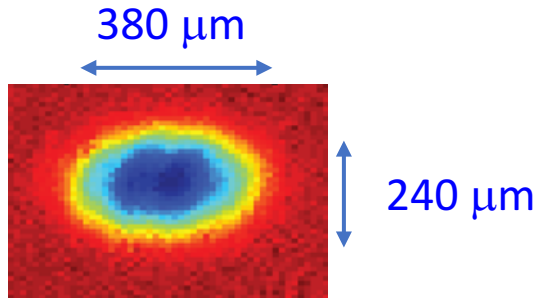
$$\delta = \omega - \omega_0$$

$$b = \text{Im}\{\alpha\} \rho k L$$

$$\phi = \frac{\text{Re}\{\alpha\}}{2} \rho k L$$

$$\alpha(\omega) = -\frac{3\pi}{k^3} \frac{\Gamma}{\delta + i\Gamma/2}$$

## Ultracold Strontium gas



Atoms number:  $N = 3 \cdot 10^8$

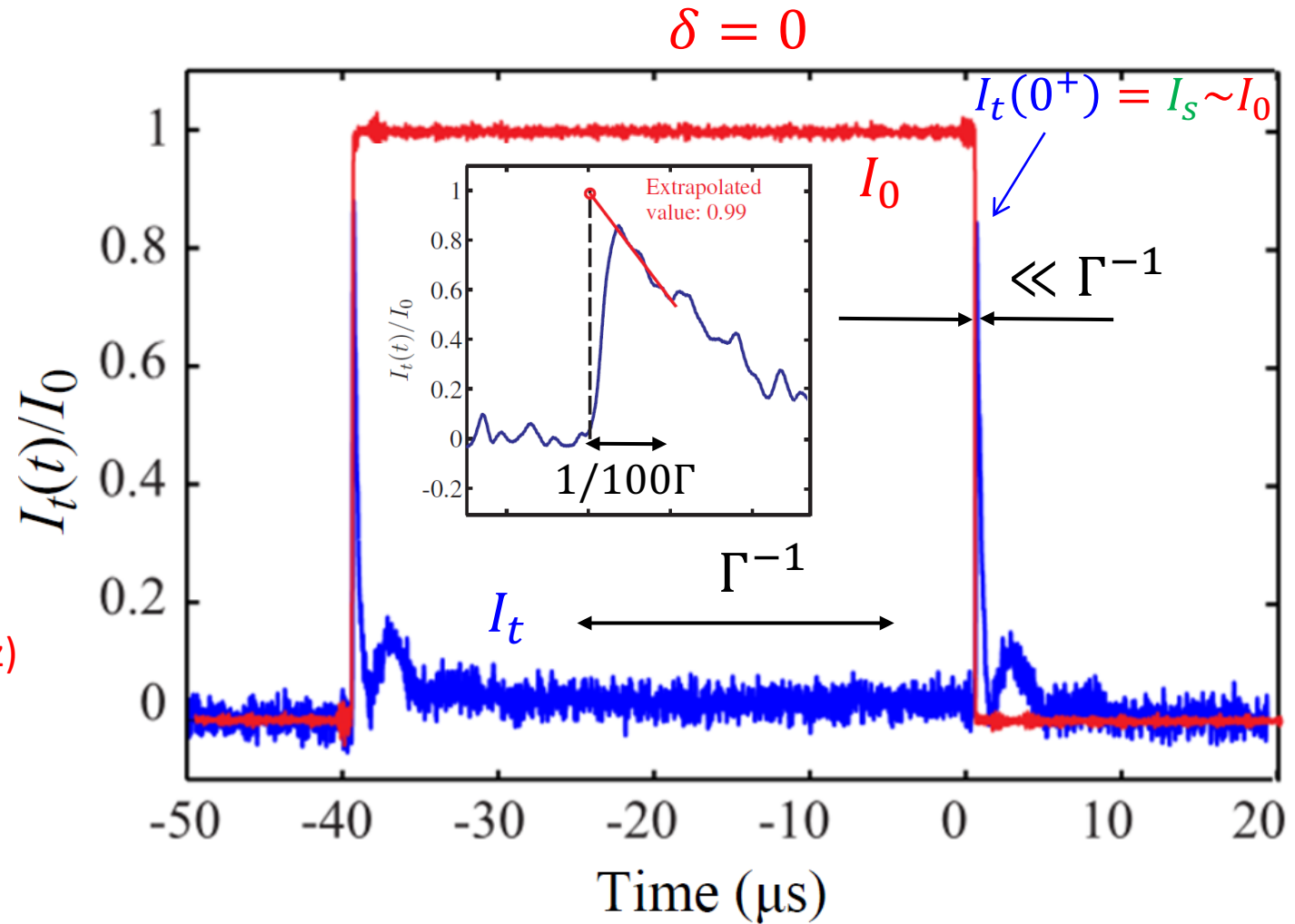
Density:  $\rho \sim 5 \cdot 10^{11}$  atoms/cm<sup>3</sup>

Dilute medium:  $\rho\lambda^3 \sim 0.1$

Temperature:  $3 \mu\text{K} \rightarrow k\bar{v} = 3.4 \Gamma$

Optical thickness:  $b_0(\bar{v}) = 19$  ( $\frac{\Gamma}{2\pi} \sim 7.5$  kHz)

$b_0(0) = 120$   
 $\delta = 0$        $T = 0$



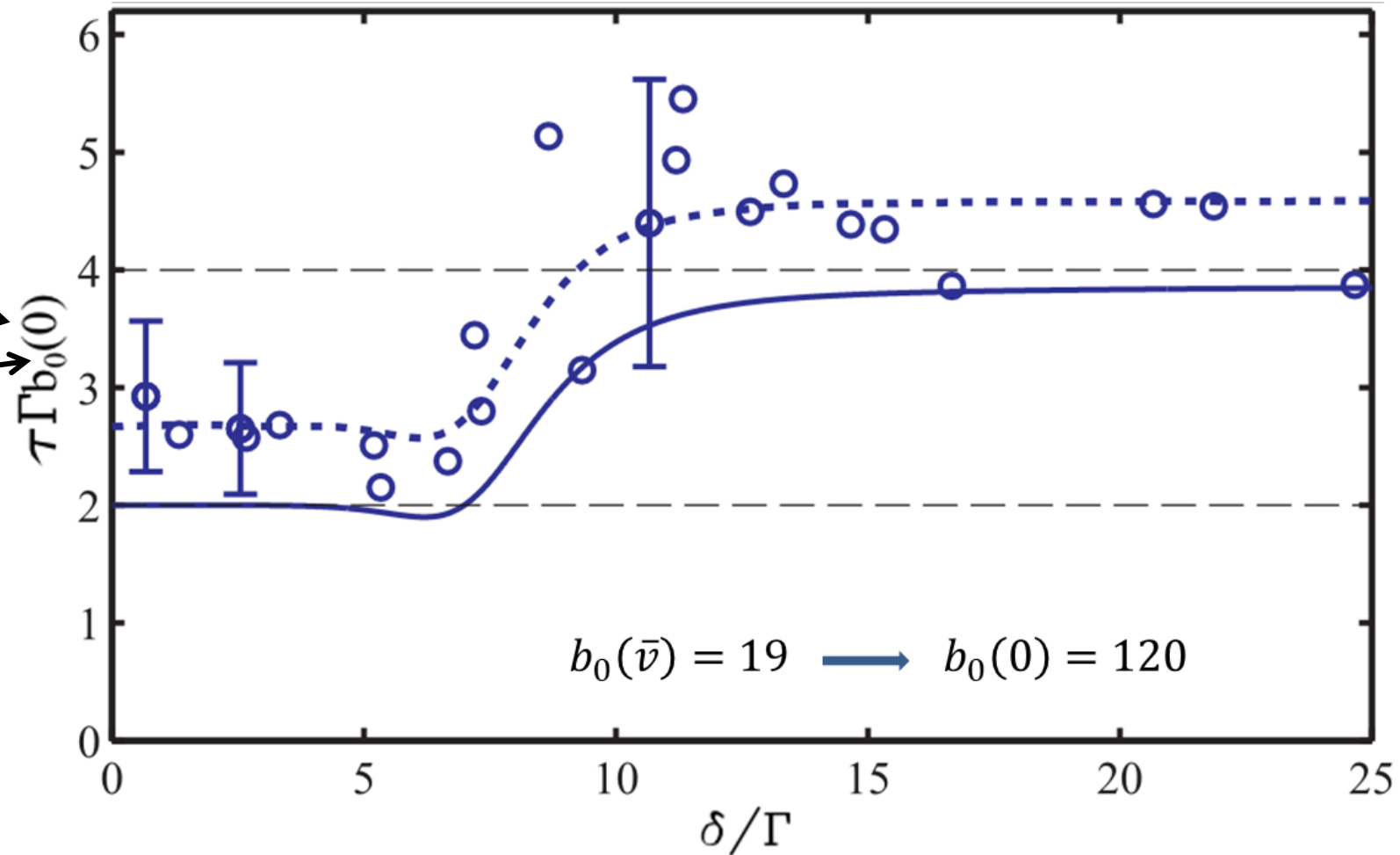
M. Chalony et al., PRA **84**, 011401(R) (2011)

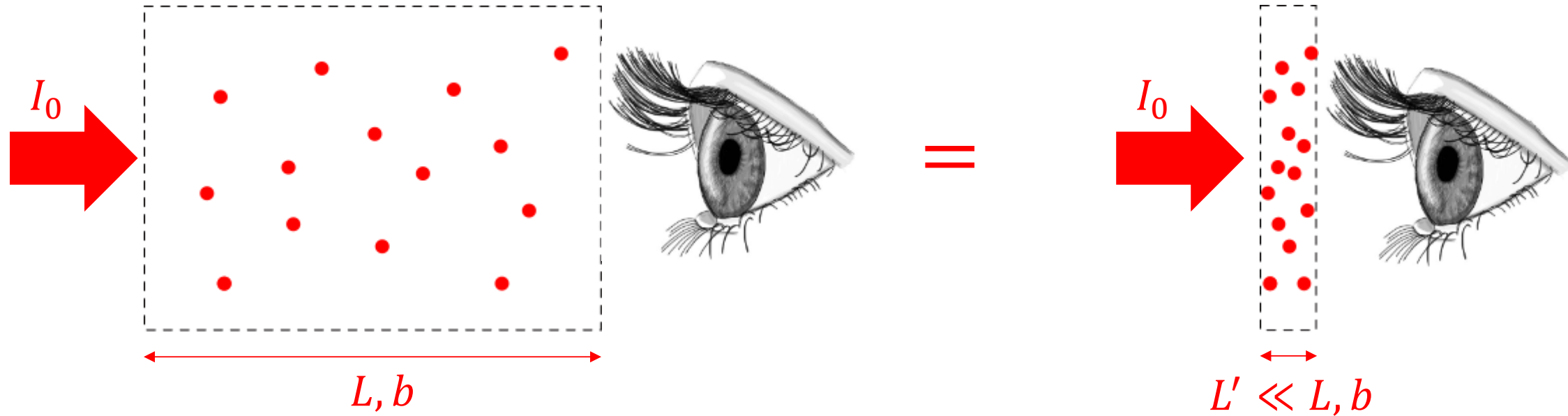
C. C. Kwong, et al. PRL **114**, 223601 (2015)

$$\tau = \frac{I(0^+)}{dI(0^+)/dt}$$

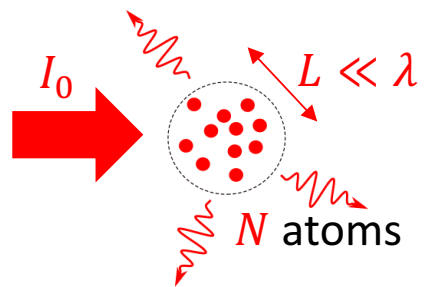
$T = 0$   
 $\delta = 0$  →  $\tau \Gamma b_0(0)$

$$\tau \sim \frac{1}{b\Gamma} \ll \frac{1}{\Gamma}$$





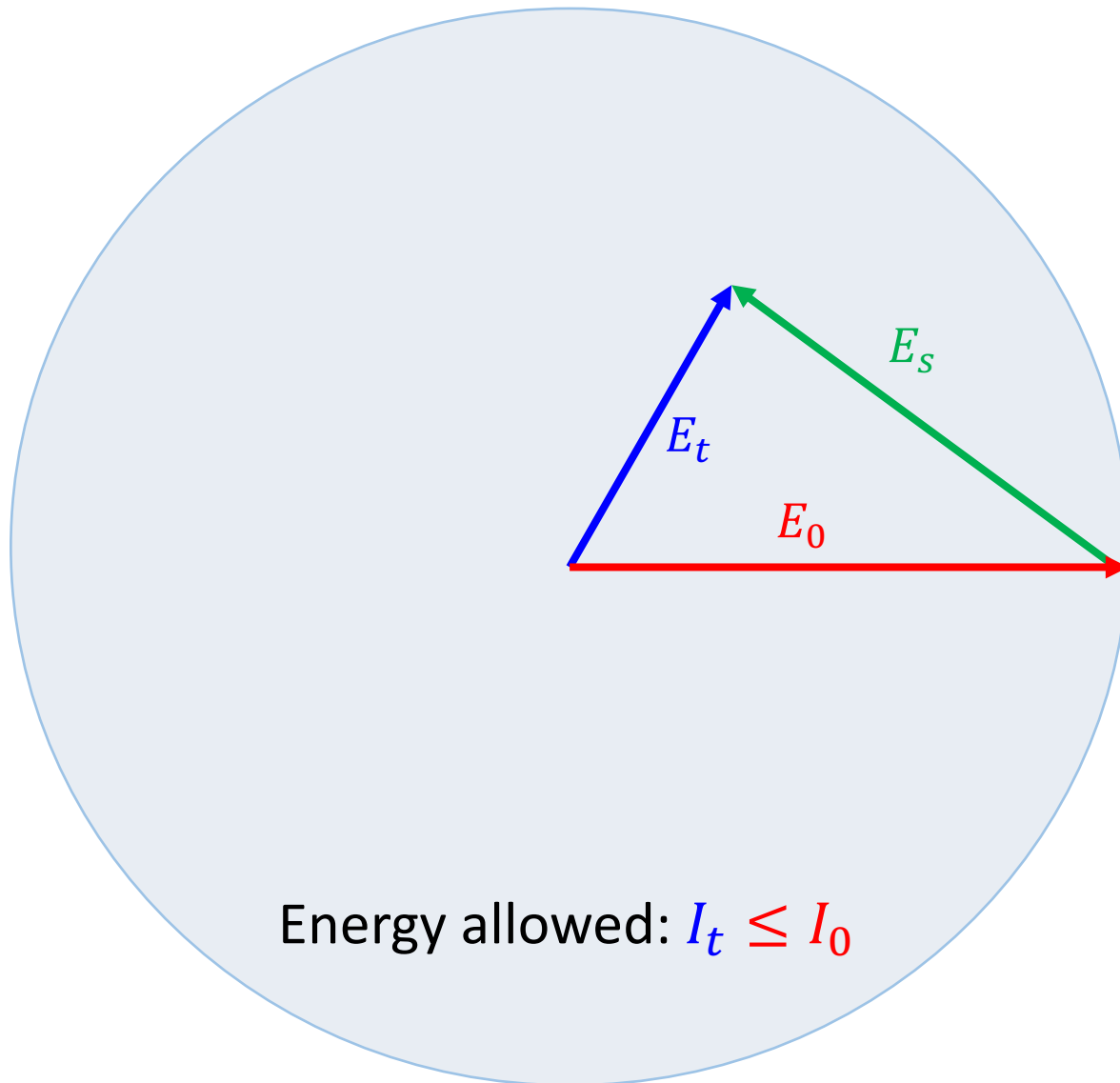
Optical thickness  $b \leftrightarrow$  Effective atoms number



Superradiance:  $\tau \sim 1/N\Gamma \leftrightarrow$  Flash  $\tau \sim 1/b\Gamma$

Flash is a cooperative emission in the forward direction



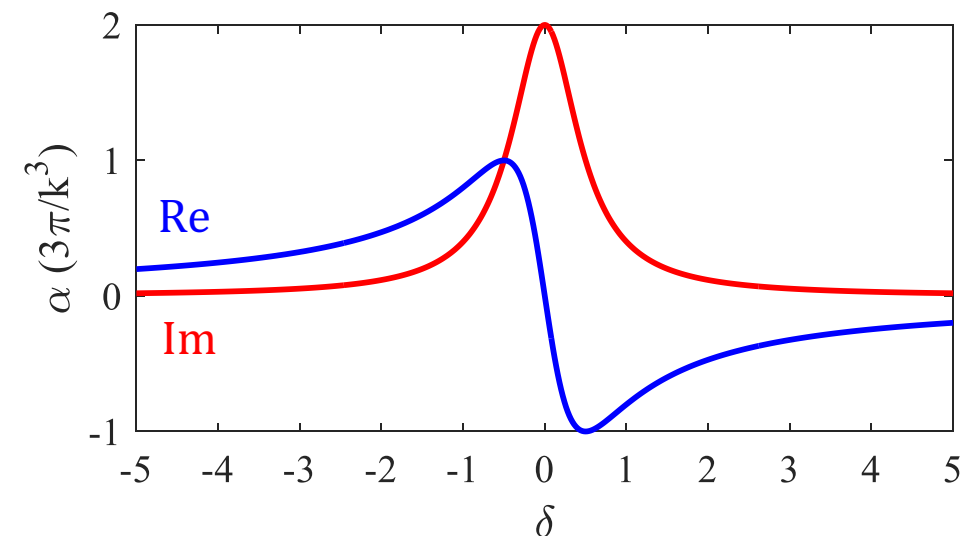


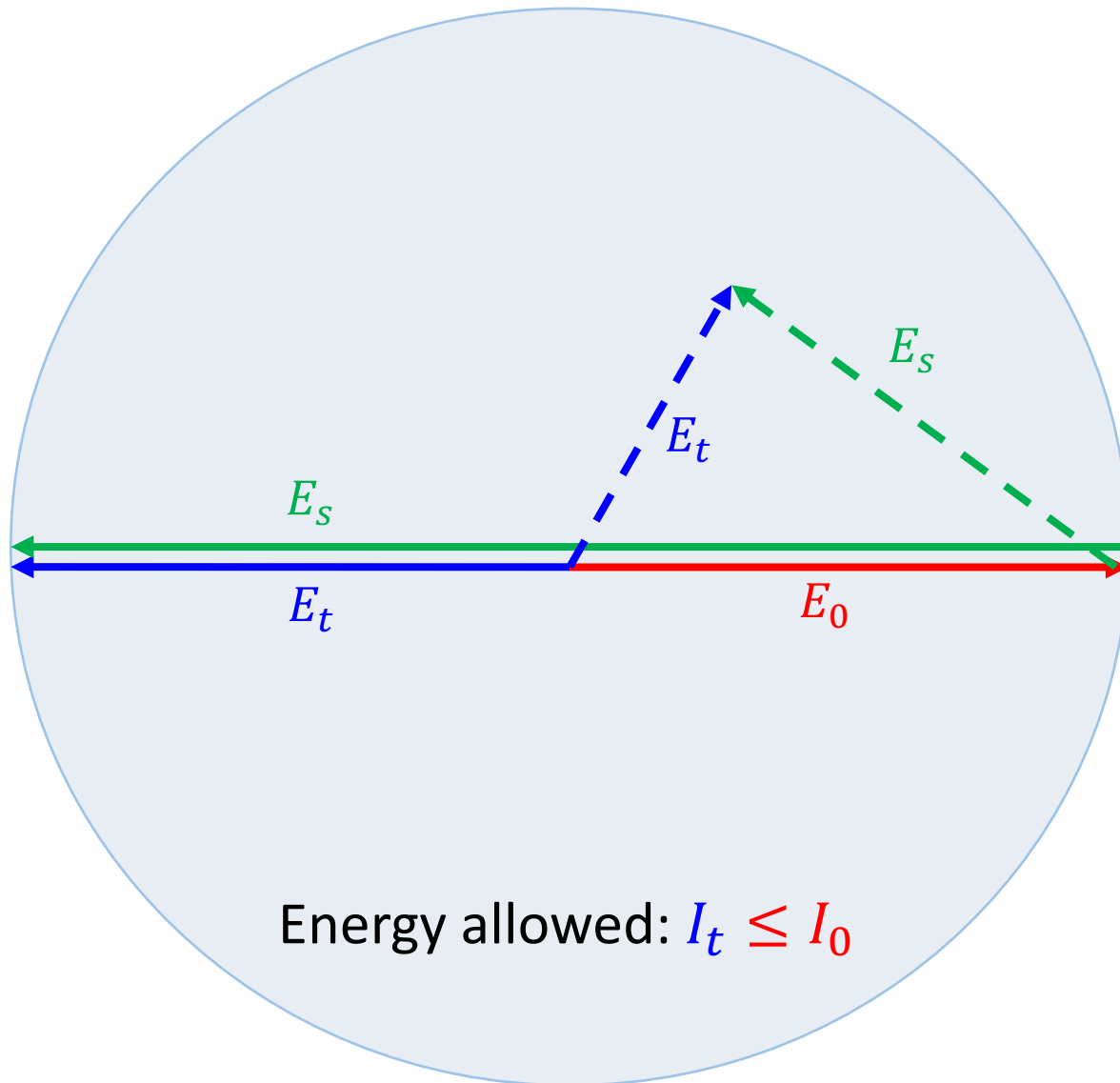
$$E_s(\omega) = E_0(\omega) \left[ \exp \left( -\frac{b}{2} + i\phi \right) - 1 \right]$$

Max. value:  $\exp \left( -\frac{b}{2} + i\phi \right) \rightarrow -1$

So we need:  $b \rightarrow 0$  and  $\phi \rightarrow \pm\pi$

$$b = \text{Im}\{\alpha\} \rho k L \quad \phi = \frac{\text{Re}\{\alpha\}}{2} \rho k L$$



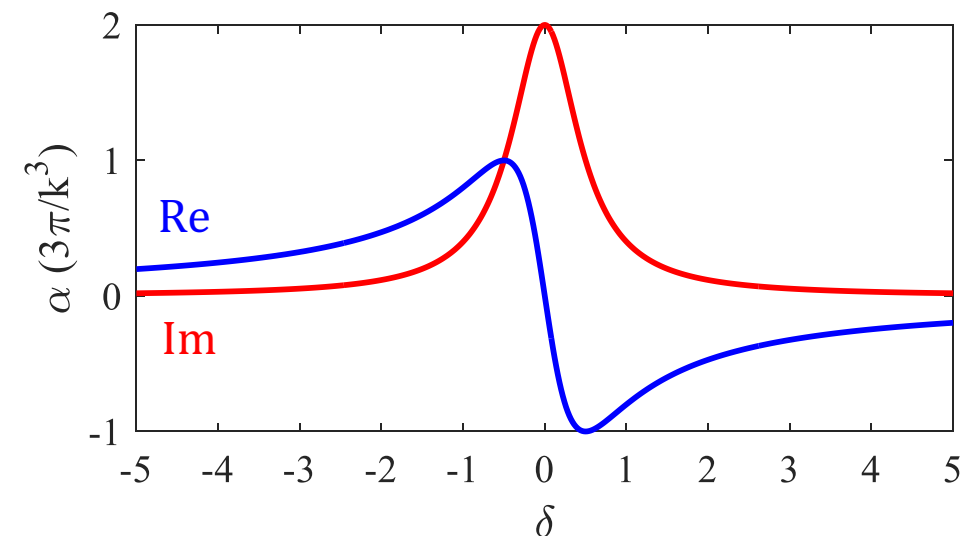


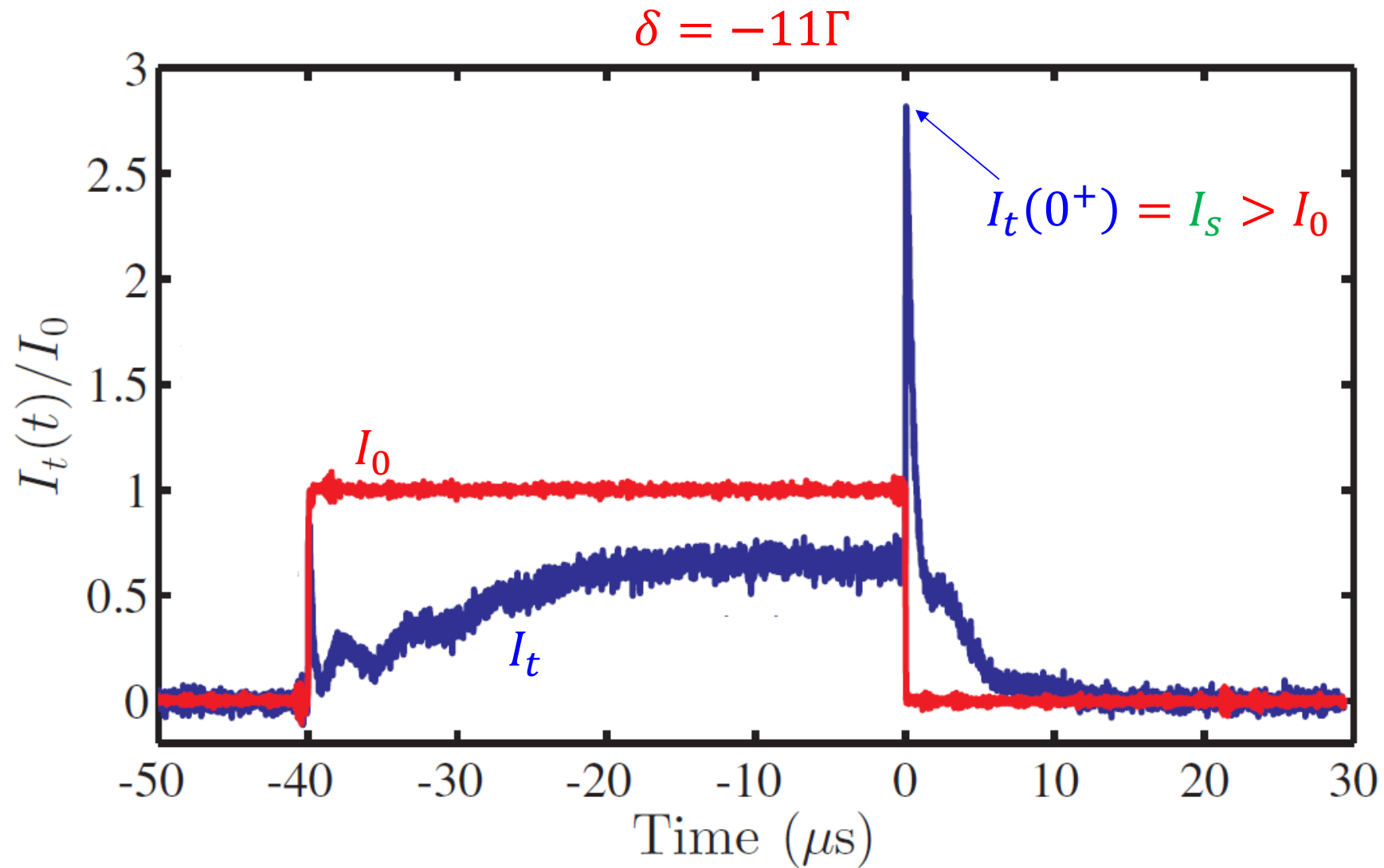
$$E_s(\omega) = E_0(\omega) \left[ \exp \left( -\frac{b}{2} + i\phi \right) - 1 \right]$$

Max. value:  $\exp \left( -\frac{b}{2} + i\phi \right) \rightarrow -1$

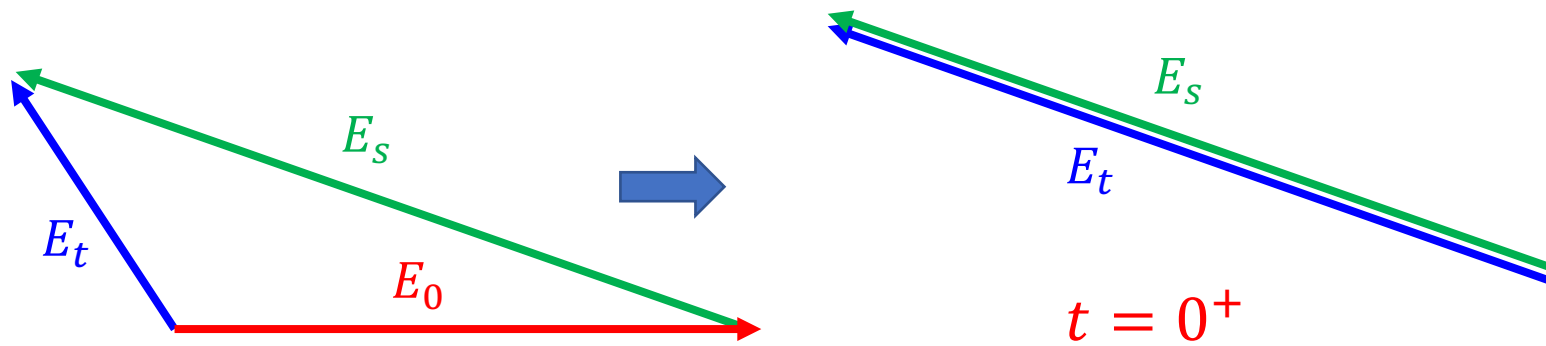
So we need:  $b \rightarrow 0$  and  $\phi \rightarrow \pm\pi$

$$b = \text{Im}\{\alpha\} \rho k L \quad \phi = \frac{\text{Re}\{\alpha\}}{2} \rho k L$$

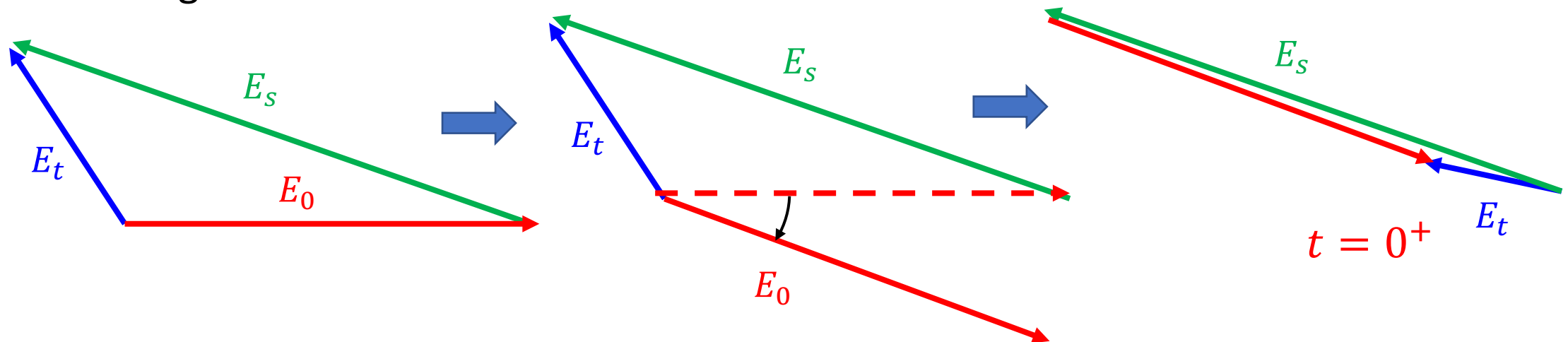




Switch off:

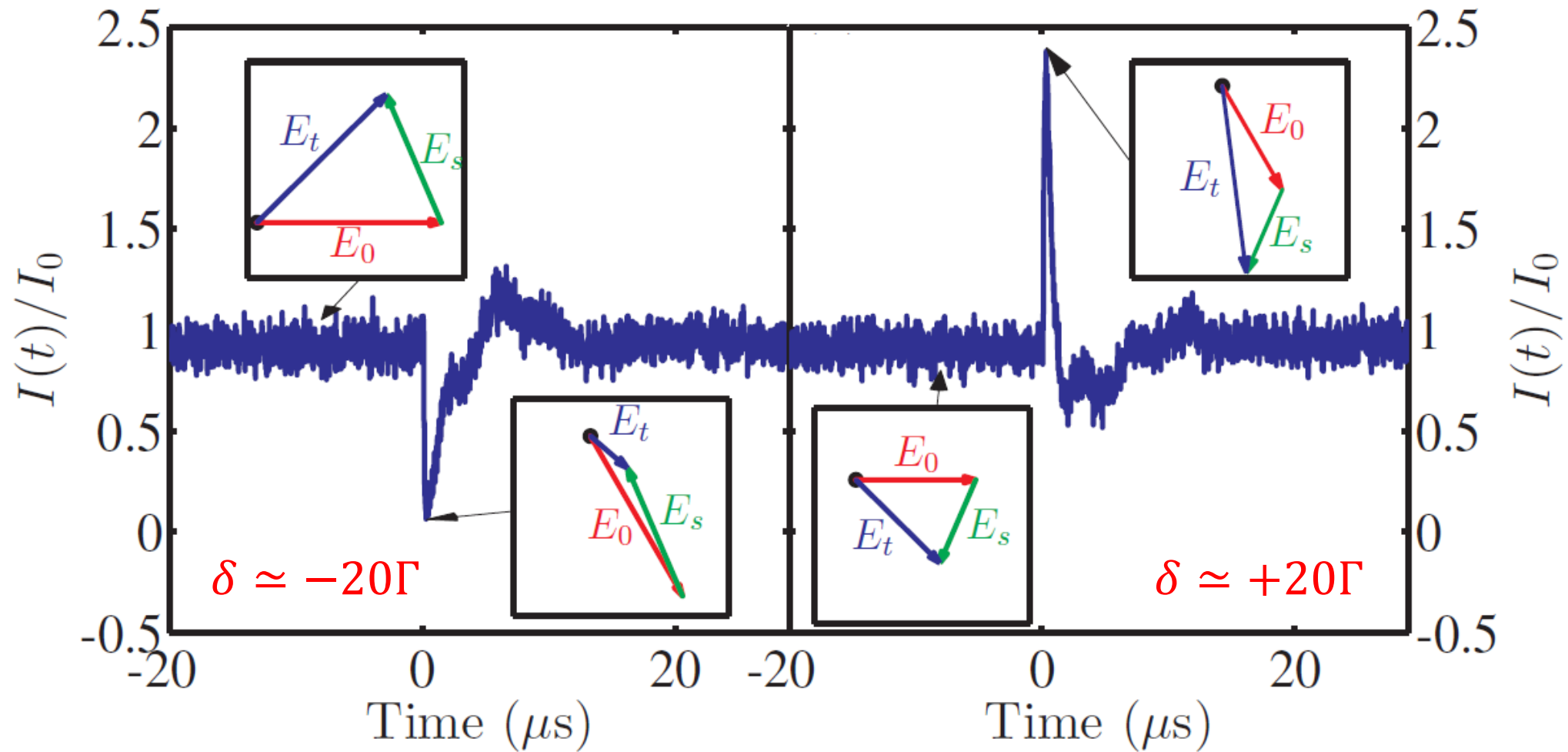


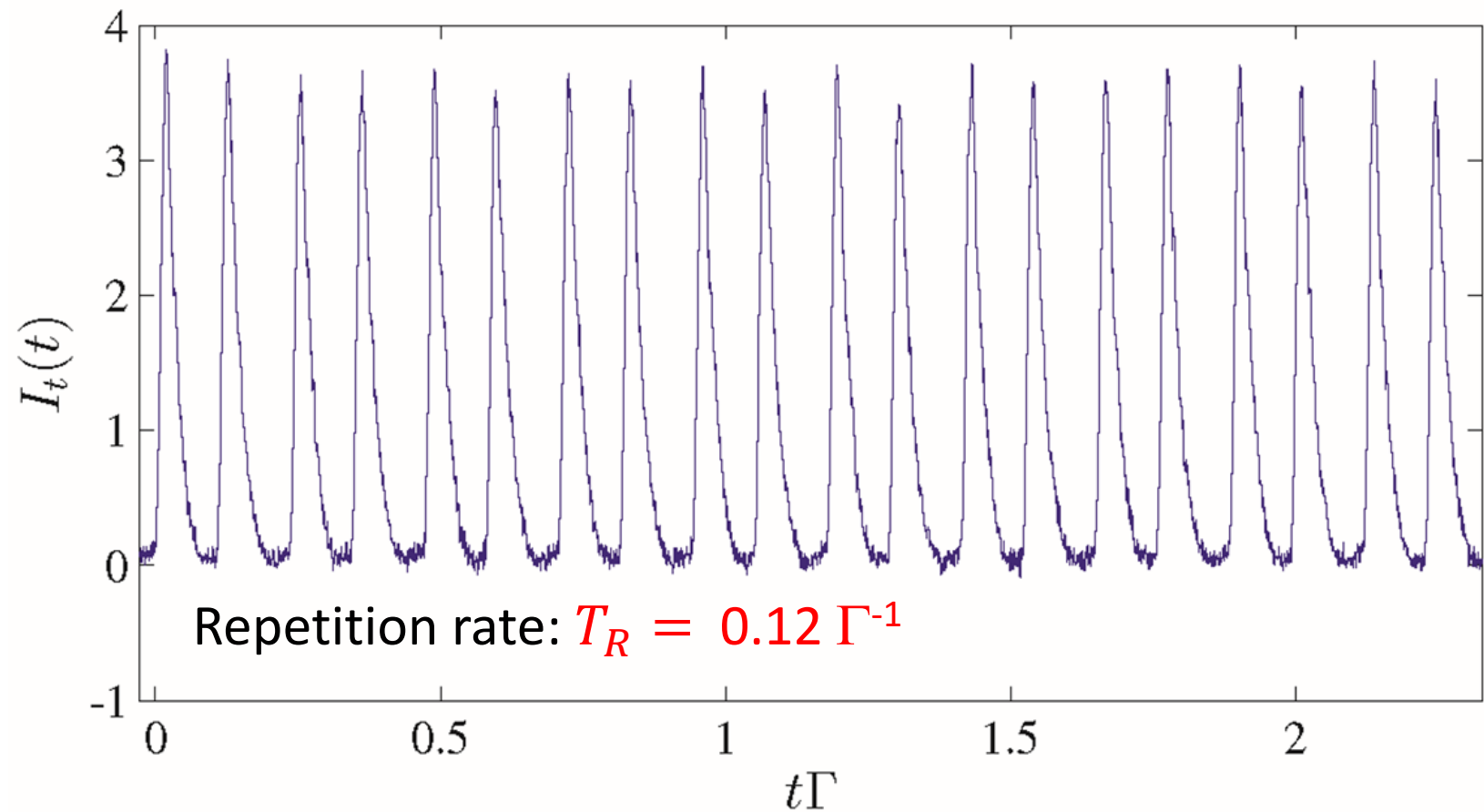
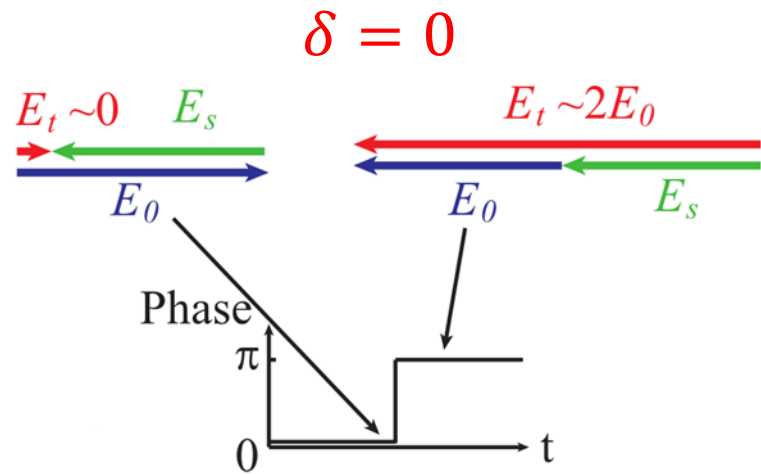
Phase change:

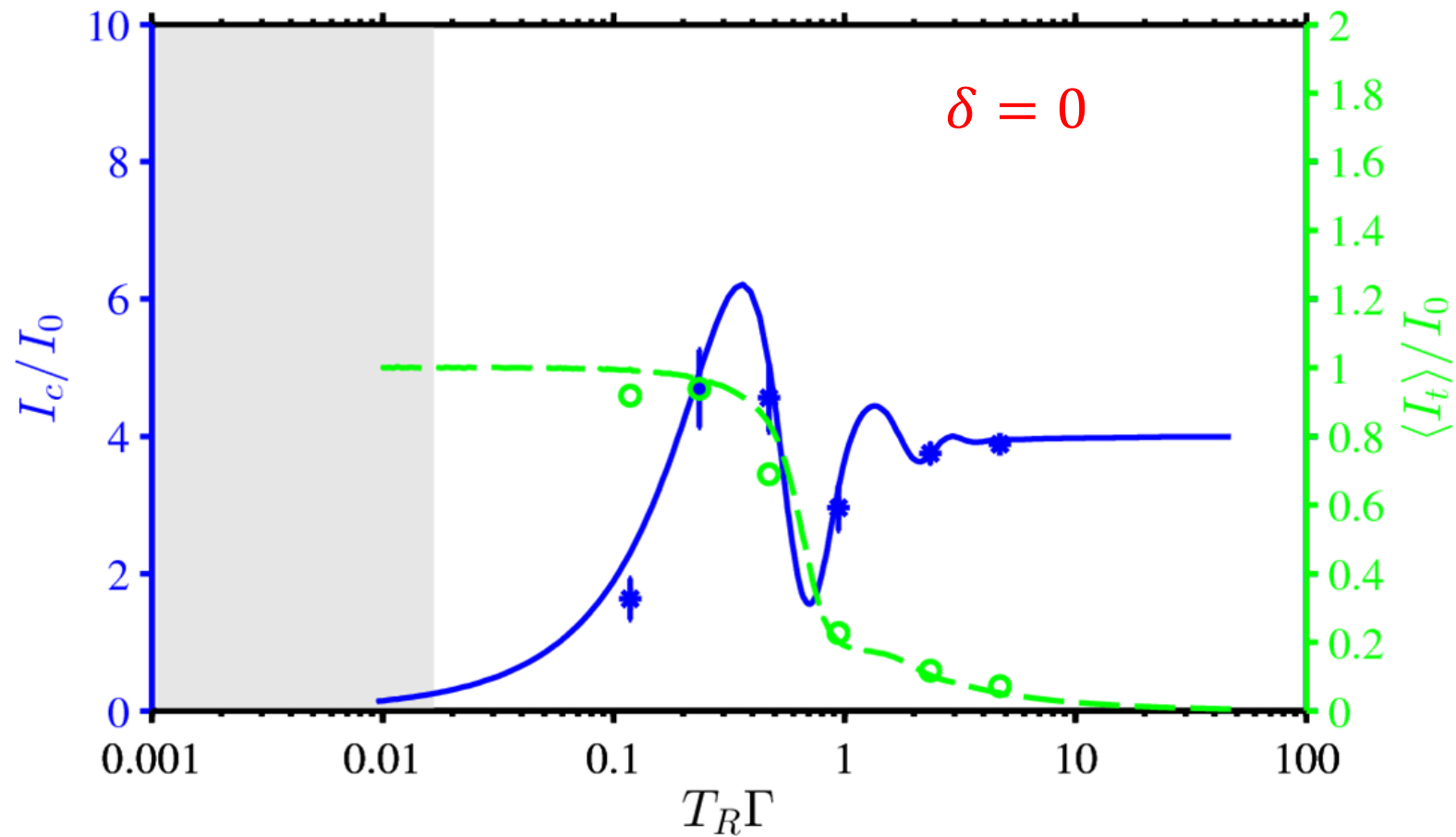




# Phase change





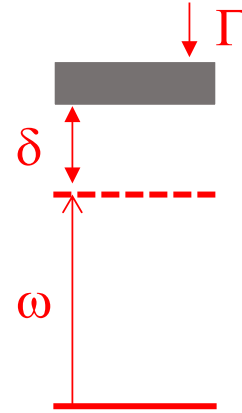


If the repetition rate  $T_R < \Gamma^{-1}$  fluorescence is quenched

Two-level:  $E_t(t) = E_0(t) + E_s(t)$

$$\alpha = -\frac{3\pi}{k^3} \frac{\Gamma}{\delta + i\Gamma/2}$$

If  $\delta = 0$  and  $b \gg 1$  ( $\phi \sim 0$ ),  $E_s \sim -E_0$



$\Lambda$ -scheme:  $E_t(t) = E_0(t) + E_s(t) + E_{S/\Lambda}(t)$

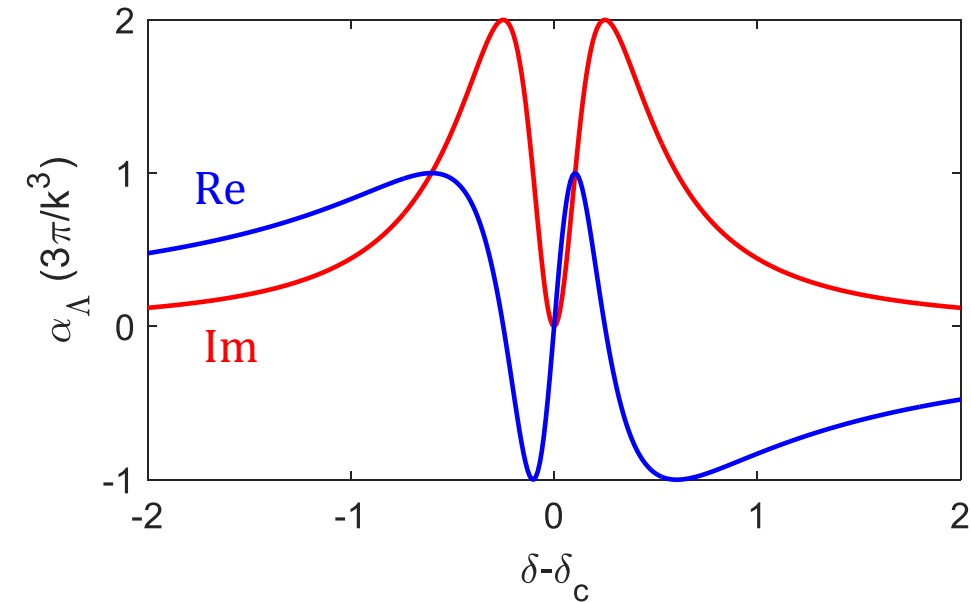
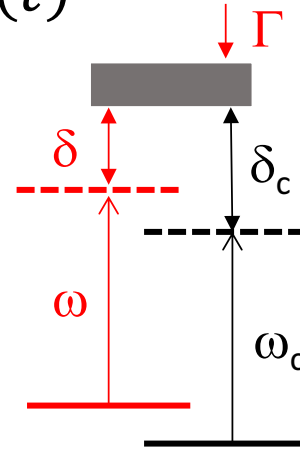
$$\alpha_\Lambda = \alpha \left( 1 - \frac{|\Omega_c|^2/4}{(\delta + i\Gamma/2)(\delta - \delta_c)} \right)^{-1}$$

Narrow dark resonance:  $\Gamma_d = |\Omega_c|^2 / \Gamma$

If  $b \gg 1$ ,  $\tau_{S/\Lambda} = b/\Gamma_d \gg \tau_S = 1/b\Gamma$

If  $\delta = \delta_c = 0$  and  $b \gg 1$  ( $\phi = 0$ ),  $E_s = -E_0$

and  $E_t = E_0$ , so  $E_{S/\Lambda} = E_0$ , and  $I_t(0^+) = |E_s + E_{S/\Lambda}|_{\text{sta}}^2 = 0$





Two-level:  $E_t(t) = E_0(t) + E_s(t)$

$$\alpha = -\frac{3\pi}{k^3} \frac{\Gamma}{\delta + i\Gamma/2}$$

If  $\delta = 0$  and  $b \gg 1$  ( $\phi \sim 0$ ),  $E_s \sim -E_0$

$\Lambda$ -scheme:  $E_t(t) = E_0(t) + E_s(t) + E_{S/\Lambda}(t)$

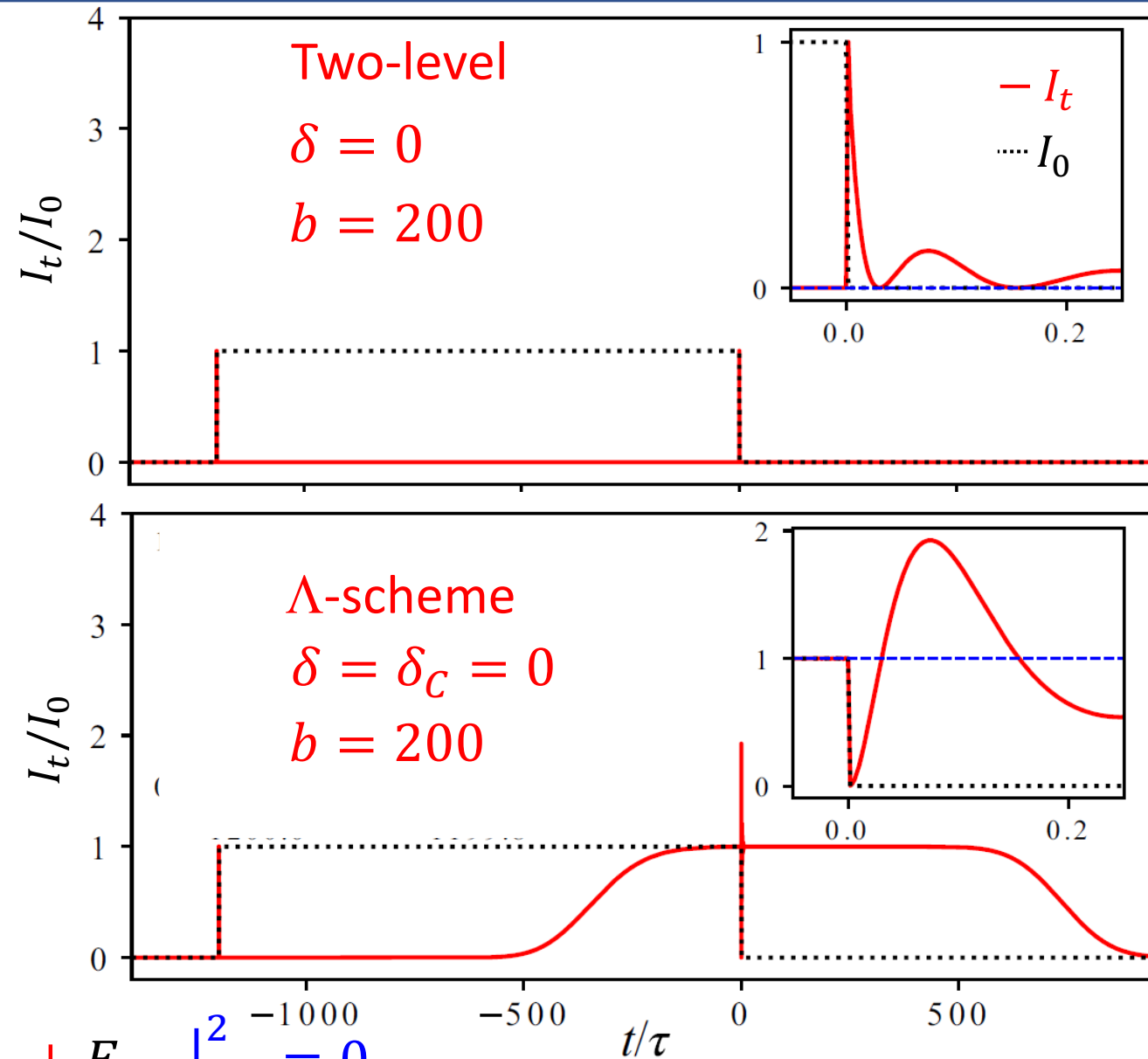
$$\alpha_\Lambda = \alpha \left( 1 - \frac{|\Omega_c|^2/4}{(\delta + i\Gamma/2)(\delta - \delta_c)} \right)^{-1}$$

Narrow dark resonance:  $\Gamma_d = |\Omega_c|^2 / \Gamma$

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If  $\delta = \delta_c = 0$  and  $b \gg 1$  ( $\phi \sim 0$ ),  $E_s = -E_0$

and  $E_t = E_0$ , so  $E_{S/\Lambda} = E_0$ , and  $I_t(0^+) = |E_s + E_{S/\Lambda}|_{\text{sta}}^2 = 0$



PRL **103**, 093602 (2009)

PHYSICAL REVIEW LETTERS

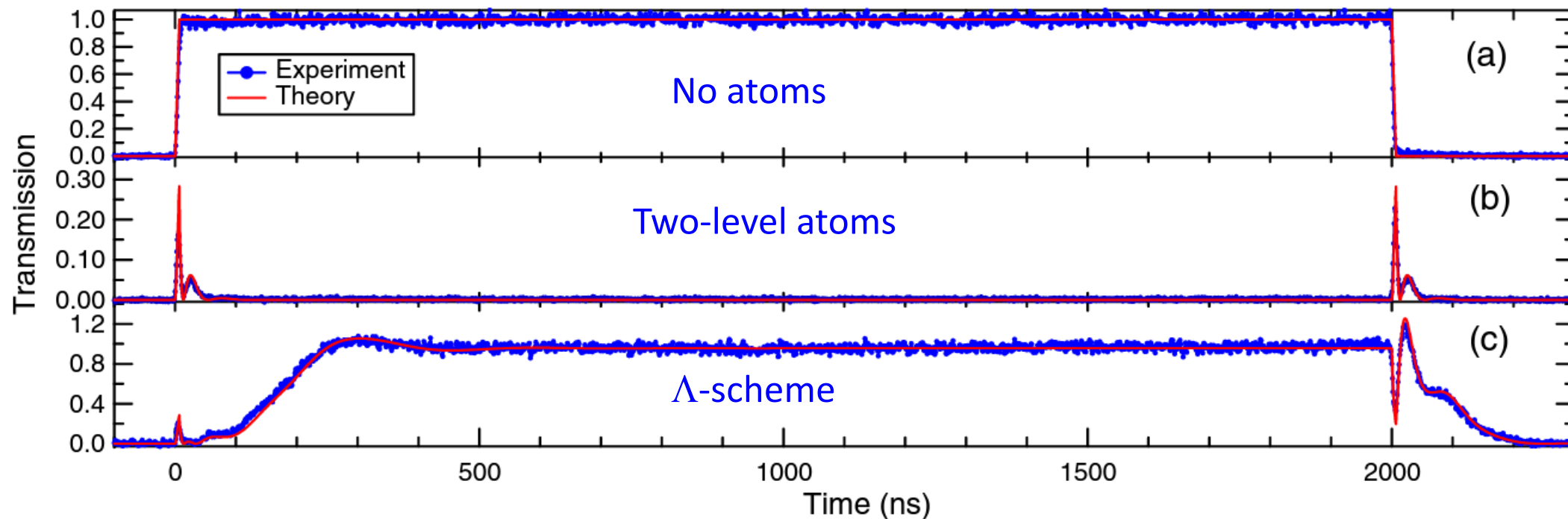
week ending  
28 AUGUST 2009

## Optical Precursors with Electromagnetically Induced Transparency in Cold Atoms

Dong Wei, J. F. Chen, M. M. T. Loy, G. K. L. Wong, and Shengwang Du\*

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(Received 10 July 2009; published 28 August 2009)

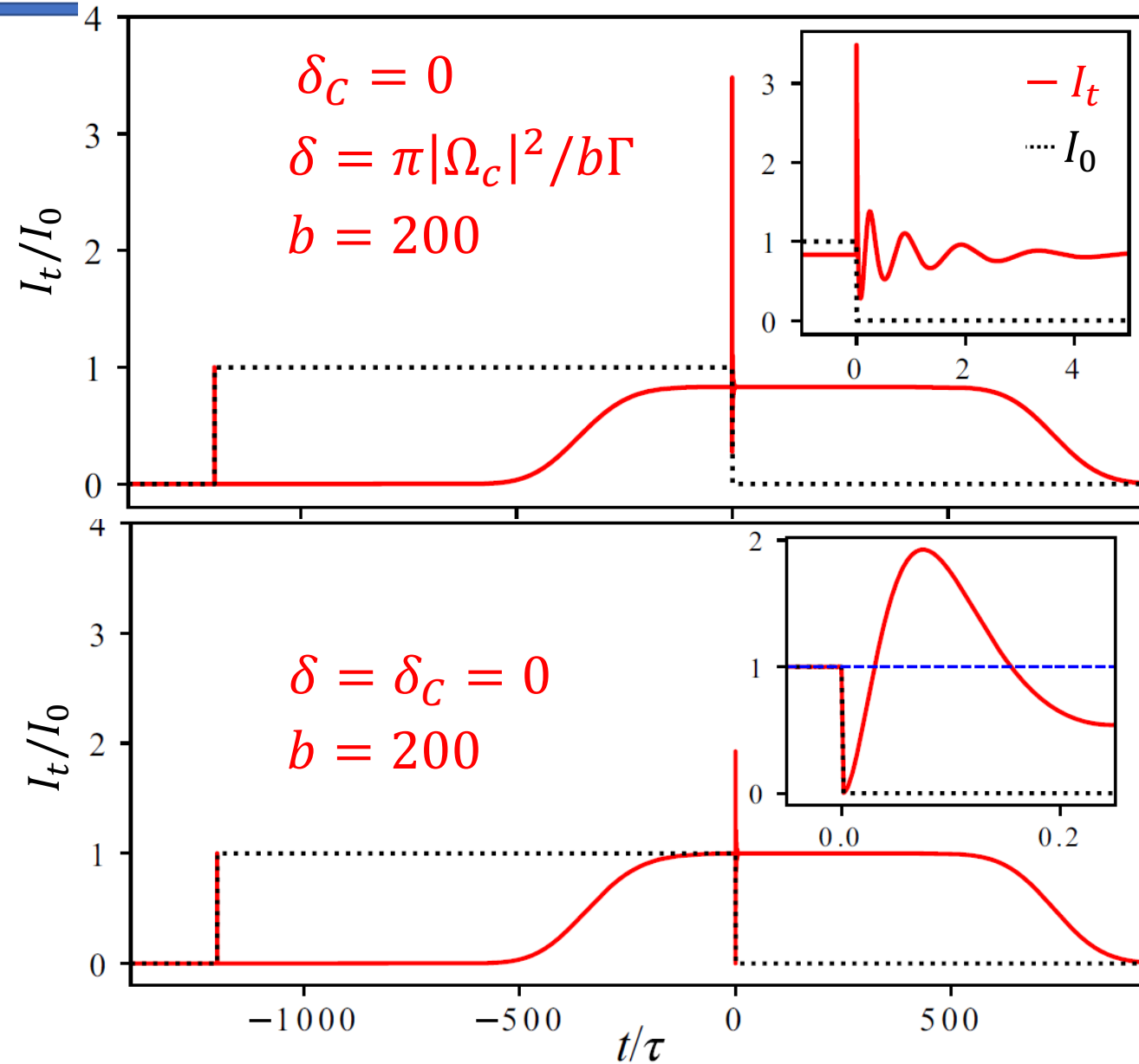
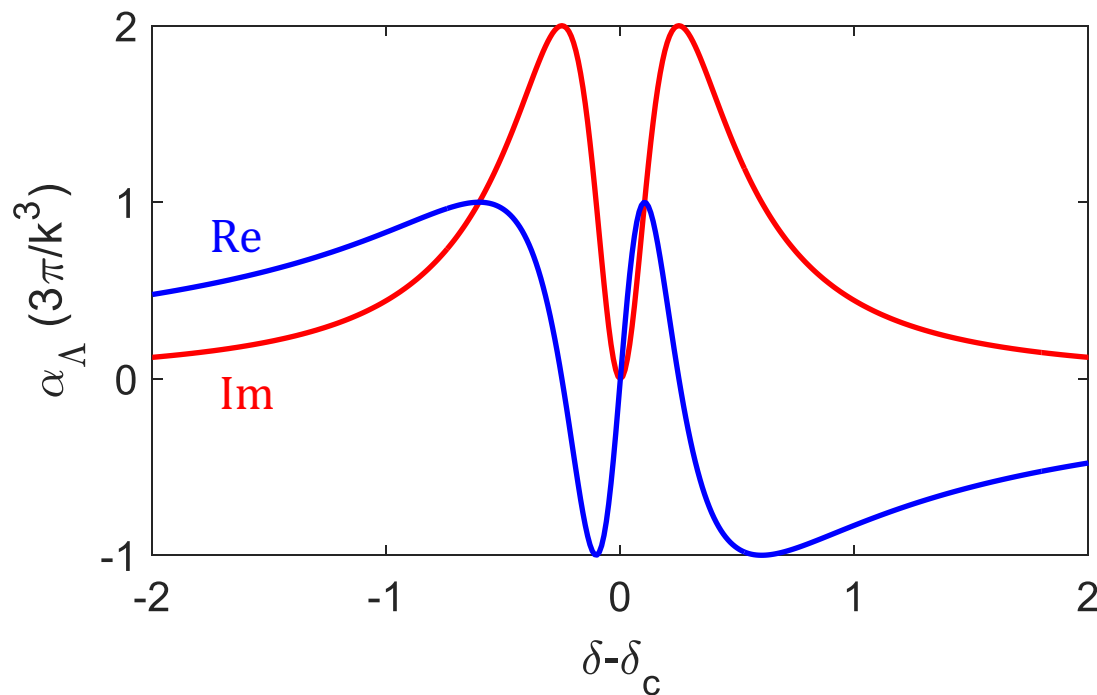


$\Lambda$ -scheme:  $E_t(t) = E_0(t) + E_s(t) + E_{S/\Lambda}(t)$

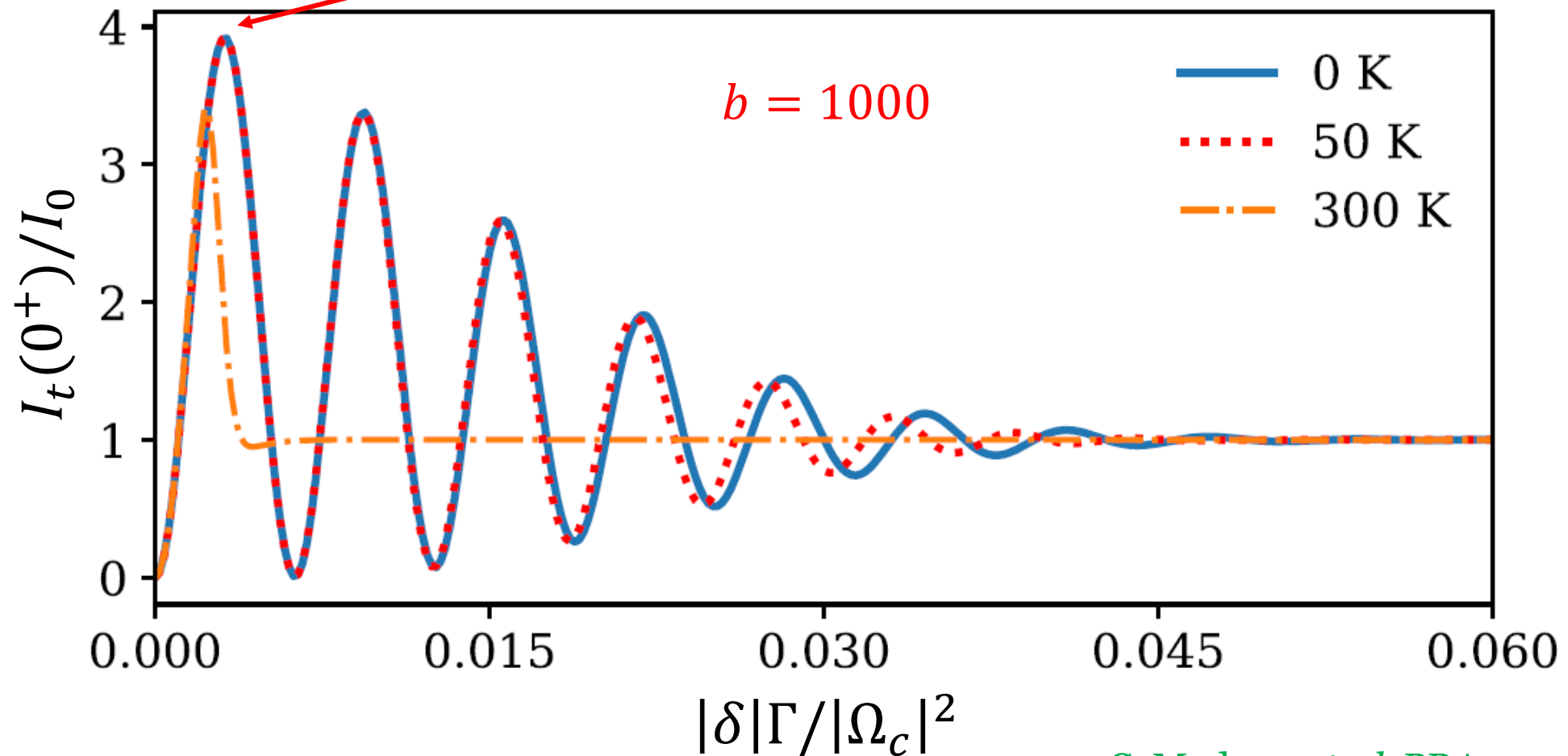
If  $\delta = \delta_c = 0$ , then  $E_{S/\Lambda} = E_0$

If  $\delta = \pm\pi|\Omega_c|^2/b\Gamma \ll \Gamma$ , then  $E_{S/\Lambda} = -E_0$

$$I_t(0^+) = |E_s + E_{S/\Lambda}|_{\text{sta}}^2 = 4$$



For Rb atoms:  $\Gamma/2\pi = 6$  MHz,  $|\Omega_c| = \Gamma/10$ ,  $|\delta|_\pi/2\pi \approx 300$  Hz and  $\frac{|\delta|_\pi}{\Delta\omega_{hf}} = 5 \times 10^{-8}$





Transient phenomena in the coherent transmission  
of a optical thick medium with “slow” resonant response time.

Decomposition of the field as:

$$E_t(t) = E_0(t) + E_s(t) + E_{S/\Lambda}(t)$$

Transmission      Incident      Forward scattering       $\Lambda$ -contribution

$E_s$  is “fast”  $\sim 1/b\Gamma$ . Similar to superradiance

$E_{S/\Lambda}$  is very “slow”  $\sim b/\Gamma$ . Slow light

$E_s$ ,  $E_{S/\Lambda}$  can be extracted using the (super)flash effect, *i.e.* measurement of  $I_t(0^+)$

Applications: Cooperative pulse generator, frequency sensing

## Flash Effect

Maryvonne Chalony

Dominique Delande

Romain Pierrat



Romain Pierrat, IL ESPCI (Fr)



## Superflash and cooperative Pulse

Kwong Chang-Chi

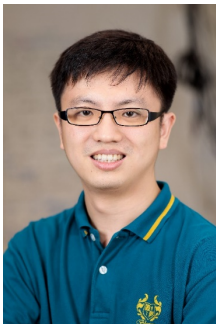
Kanhaiya Pandey

Mysore Pramod

Tao Yang

Dominique Delande

Romain Pierrat



M. Chalony et *al.*, PRA **84**, 011401(R) (2011)

Dominique Delande, LKB SU (Fr)



C. C. Kwong, et *al.*, PRL **113**, 223601 (2014)

C. C. Kwong, et *al.* PRL **114**, 223601 (2015)

## Flash with $\Lambda$ -Scheme

Chetan S. Madasu

Kwong Chang-Chi

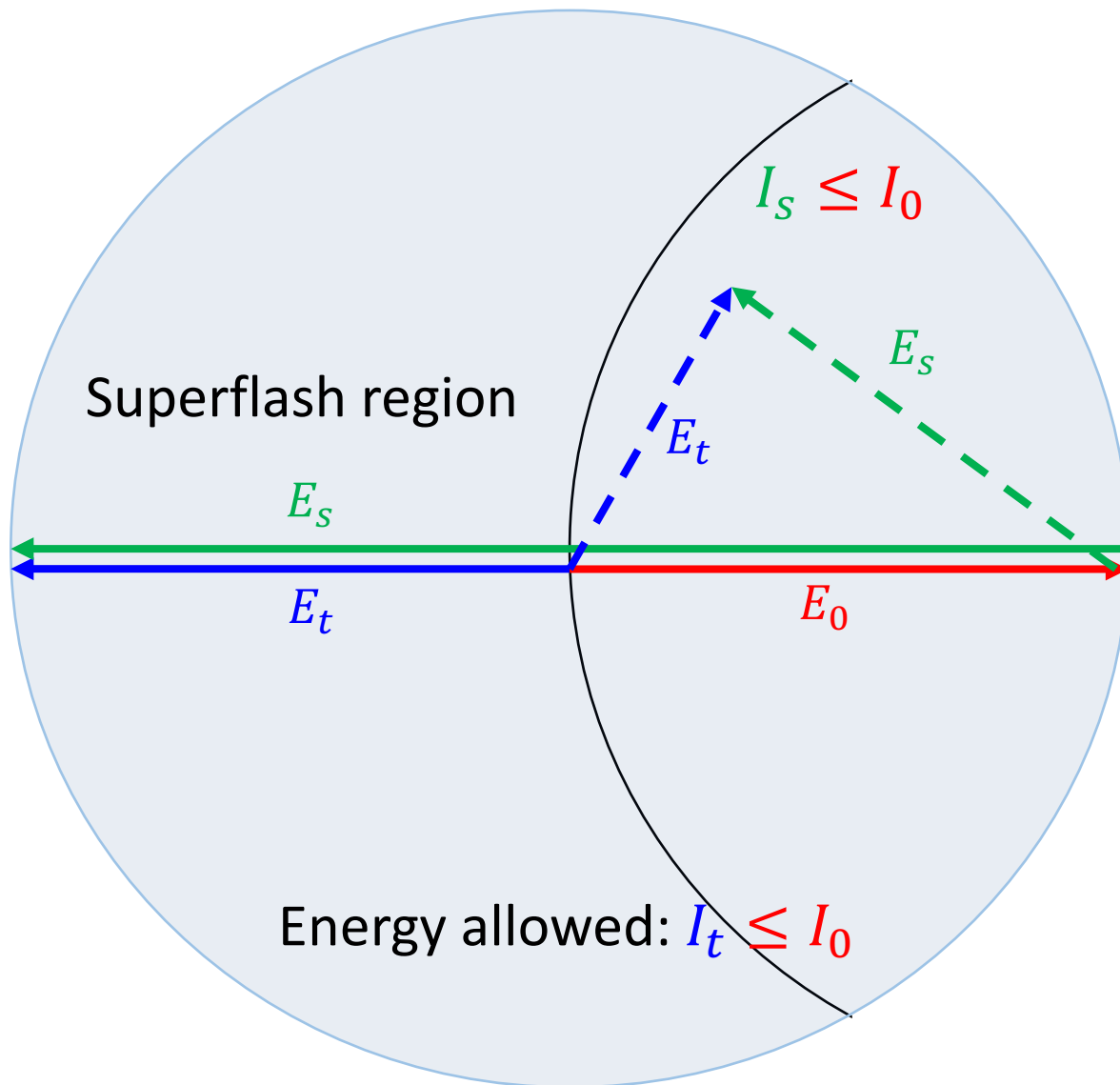
Kanhaiya Pandey



C. Madasu, et *al.* PRA **105**, 013713 (2022)

Kanhaiya Pandey, IIT Guwahati (In)





$$E_s(\omega) = E_0(\omega) \left[ \exp \left( -\frac{b}{2} + i\phi \right) - 1 \right]$$

Max. value:  $\exp \left( -\frac{b}{2} + i\phi \right) \rightarrow -1$

So we need:  $b \rightarrow 0$  and  $\phi \rightarrow \pm\pi$

$$b = \text{Im}\{\alpha\} \rho k L \quad \phi = \frac{\text{Re}\{\alpha\}}{2} \rho k L$$

