

Scattering of topological edge waves in Kekule structures

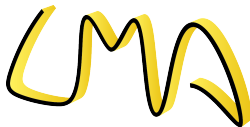
Antonin Coutant

8 Dec. 2022

GdR Complexe: Annual workshop

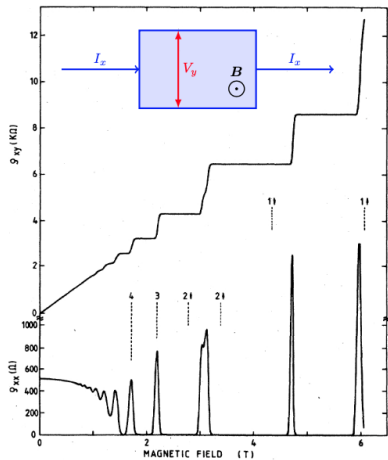
Collaboration with:

V. Achilleos, V. Pagneux, O. Richoux, G. Theocharis, L.-Y. Zheng



Laboratoire de Mécanique et d'Acoustique

Topological insulators: origins



First discovery:

- Quantum Hall effect (1986)

Later \rightarrow broad field:

- In photonics (\sim 2008), in acoustics or elasticity (\sim 2015)

Interest:

Topological protection

- Properties **unchanged** through **continuous transformations**

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- 2D: Transport robust against defects or disorder
→ **Perfect transmission**

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Topological protection

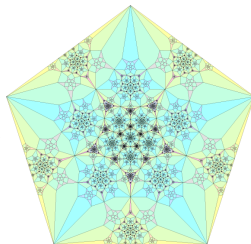
- Properties **unchanged** through **continuous transformations**
- 2D: Transport robust against defects or disorder
→ **Perfect transmission?**

Interest:**Topological protection**

- Properties **unchanged** through **continuous transformations**
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Topology



Symmetry

- **Careful:** topology is **not** enough
- **Need also** symmetries (broken or maintained)

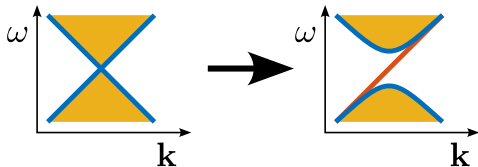
Outline

- 1 2D topological waves
 - Quantum Hall Effect
 - Quantum Spin Hall Effect
- 2 The third way: chiral and mirror symmetries
 - Chirality and SSH model
 - Kekule model
 - Mirror winding numbers
- 3 Chiral-Mirror edge mode scattering
 - Symmetries
 - Scattering results

Topological modes with classical waves:

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Quantum Hall effect:

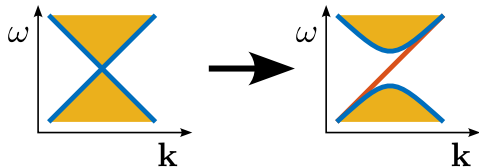


Symmetry class: **Broken time reversal**

Topological invariant: Chern number

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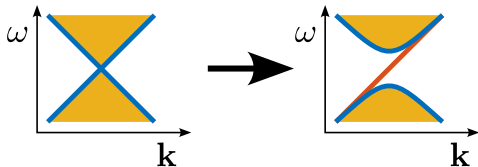
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- Backscattering immunity:
guaranteed by **unidirectional mode**

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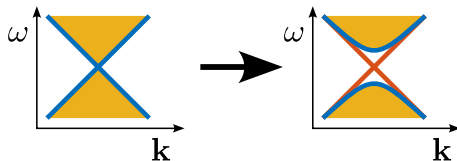
Symmetry class: **Broken time reversal**

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- Backscattering immunity:
guaranteed by **unidirectional mode**
- Drawbacks:
 - Active materials \rightarrow energy to maintain TR breaking
 - Ex. in **acoustic**: Flow \rightarrow dissipation/instabilities

Topological modes with classical waves:

Quantum Spin Hall effect:

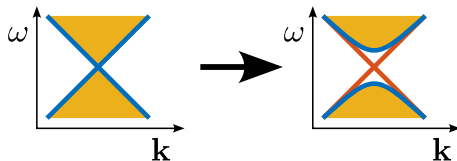


Symmetry class: **Time reversal and spin**

Topological invariant: \mathbb{Z}_2 invariant

Topological modes with classical waves:

Quantum Spin Hall effect:



Symmetry class: **Time reversal and spin**

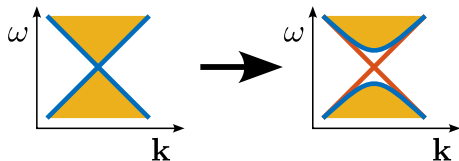
Topological invariant: \mathbb{Z}_2 invariant

- Backscattering immunity: guaranteed by

$$\mathcal{T}^2 = -1 \quad (\text{Time reversal operator})$$

Topological modes with classical waves:

Quantum Spin Hall effect:



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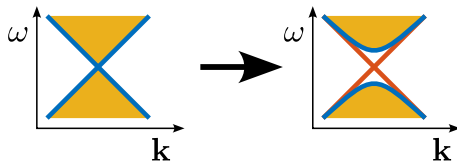
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- Drawbacks:

- Classical waves **always** have $\mathcal{T}^2 = 1$ (Bosons)

Topological modes with classical waves:

Quantum Spin Hall effect:



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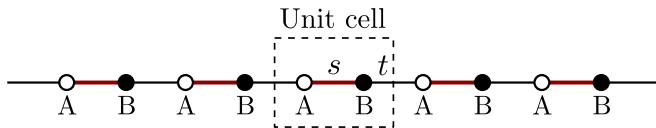
- Drawbacks:

- Classical waves **always** have $\mathcal{T}^2 = 1$ (Bosons)
- Can mimick a spin with $\mathcal{T}_{\text{eff}}^2 = -1$
 - Spatial symmetries (but broken by defects)
 - Valley degrees of freedom (but always approximate)

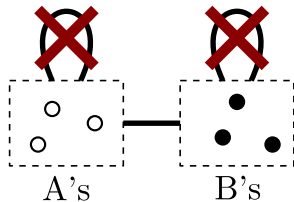
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Chiral symmetry



Ex: SSH model has chiral symmetry



- Chiral \Leftrightarrow **sublattice** symmetry
- **Algebra:** unitary matrix Γ

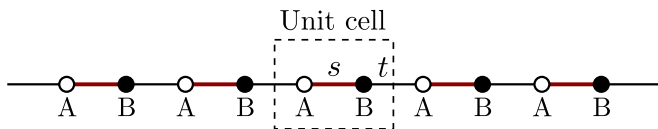
$$\Gamma H + H \Gamma = 0$$

with H Hamiltonian

- First consequence: spectrum symmetry

$$\varepsilon \rightarrow -\varepsilon$$

Chiral symmetry



Ex: SSH model has chiral symmetry

- 1D chiral systems (class AIII)
- Topological invariant: winding number $\nu \in \mathbb{Z}$

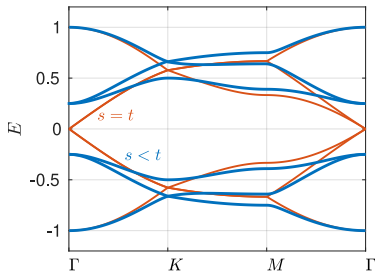
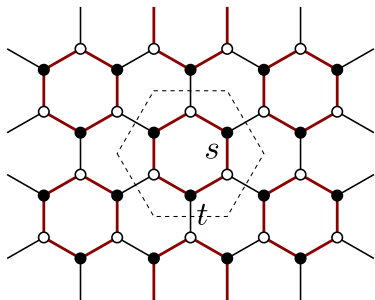
$$H(q) = \begin{pmatrix} 0 & Q \\ Q^\dagger & 0 \end{pmatrix} \quad \text{defines} \quad \nu = \frac{1}{2\pi} \int \text{Tr}(Q^{-1} \partial_q Q)$$

- $\nu \neq 0$: **Topologically protected edge mode at $\varepsilon = 0$**

In 2D, need something more:

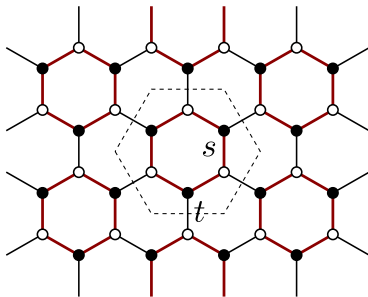
Combination of symmetries

Kekule model

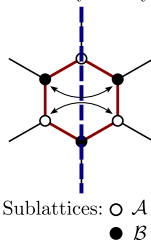


- Graphene with hopping modulation
- Intra-molecular hopping $s \neq$ extra-molecular hopping t
- Usually seen as a Quantum Spin Hall analogue

Kekule model

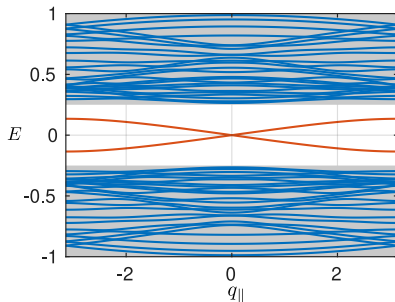
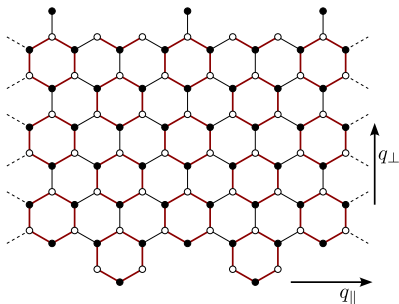


Mirror symmetry



- Chiral symmetry
- Spatial symmetries: hexagonal group C_{6v}
- **In particular:** mirror symmetries along summits $M_{j=1..3}$

$$\Gamma M_j = M_j \Gamma$$

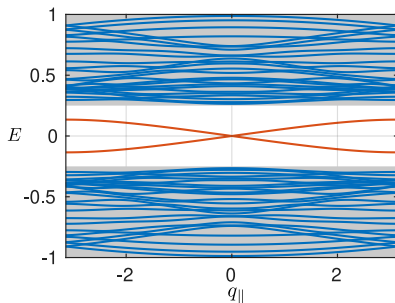
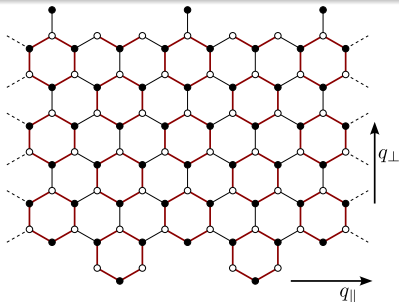


Idea of mirror winding number:

- $q = 0$ (high symmetry point):
 H splits in Symmetric and Anti-symmetric parts

$$H_S(q_\perp) = \begin{pmatrix} 0 & Q_S \\ Q_S^\dagger & 0 \end{pmatrix} \quad \text{and} \quad H_A(q_\perp) = \begin{pmatrix} 0 & Q_A \\ Q_A^\dagger & 0 \end{pmatrix}$$

- Each part is chiral \rightarrow 1D winding number



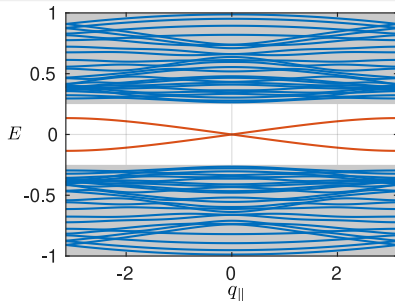
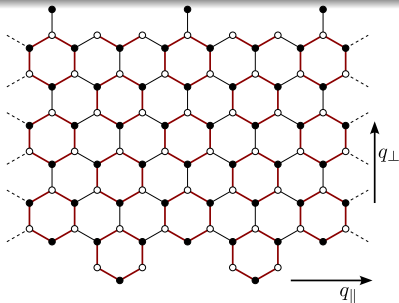
Idea of mirror winding number:

- We obtain a pair of mirror winding number:

$$\text{Topological: } (n_S, n_A) = (1, -1)$$

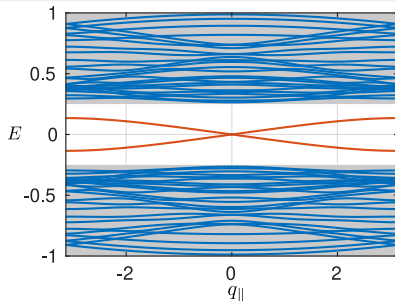
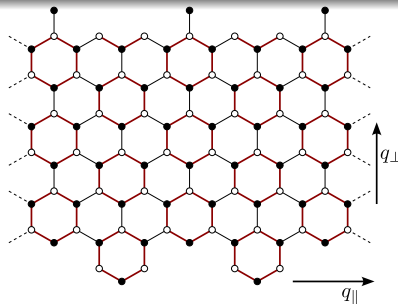
$$\text{Trivial: } (n_S, n_A) = (0, 0)$$

- 2 zero energy modes at $\varepsilon = 0$
→ **Edge waves with Dirac point**



Careful:

- Choice of edge shape matters!
- Choose edge **mirror invariant**



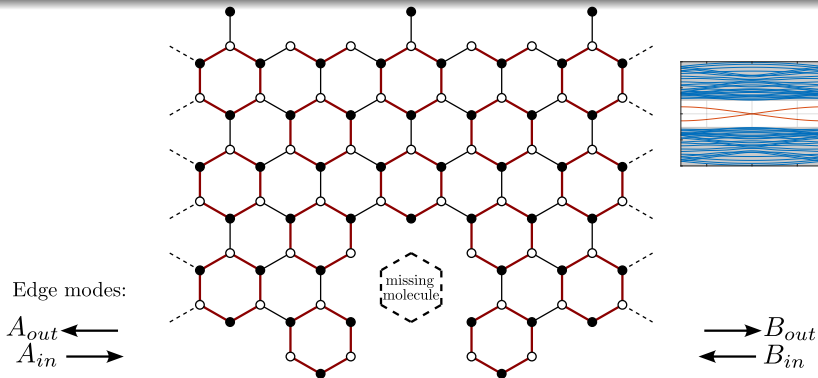
Careful:

- Choice of edge shape matters!
 - Choose edge **mirror invariant**
 - Finite Ribbon:
 - Down: molecular zigzag \rightarrow topological when $s < t$
 - Up: Partially bearded \rightarrow topological when $s > t$
- \rightarrow **exact crossing for finite width** (no mini-gap)

[AC, Zheng, Achilleos, Richoux, Theocharis, Pagneux *to appear*]

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- Scattering problem:

$$\begin{pmatrix} A_{out} \\ B_{out} \end{pmatrix} = \begin{pmatrix} R(\varepsilon) & \tilde{T}(\varepsilon) \\ T(\varepsilon) & \tilde{R}(\varepsilon) \end{pmatrix} \begin{pmatrix} A_{in} \\ B_{in} \end{pmatrix}$$

- We consider **mirror symmetric** and **chiral** defects

[AC, Zheng, Achilleos, Richoux, Theocharis, Pagneux *to appear*]

Consequences of symmetries

- S -matrix:

$$S(\varepsilon) = \begin{pmatrix} R & \tilde{T} \\ T & \tilde{R} \end{pmatrix}$$

- **Unitarity** (energy conservation)

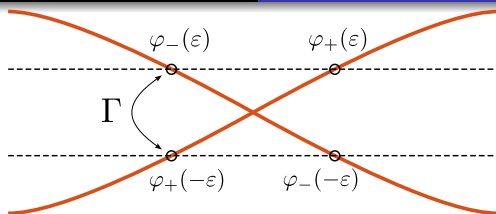
$$\begin{aligned} |R|^2 + |T|^2 &= 1, \\ R\tilde{T}^* + T\tilde{R}^* &= 0, \end{aligned}$$

- **Mirror symmetry**

$$R(\varepsilon) = \tilde{R}(\varepsilon) \quad \text{and} \quad T(\varepsilon) = \tilde{T}(\varepsilon)$$

- **Both:**

$$\text{Re}(R(\varepsilon)T^*(\varepsilon)) = 0$$



Consequences of symmetries

- **Chiral symmetry:** relation between S and $S^{-1} = S^\dagger$
- Identity:

$$R(-\varepsilon) = R(\varepsilon)^* \quad \text{and} \quad T(-\varepsilon) = T(\varepsilon)^*$$

- At $\varepsilon = 0$: real-valued
- Combining all symmetries:

$$R(0)T(0) = 0$$

Conclusion:

- Combination of symmetries imposes:

$$R(0) = 0 \quad \text{or} \quad T(0) = 0$$

Perfect transmission

Perfect reflection

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- When is it **perfect transmission**?

Conclusion:

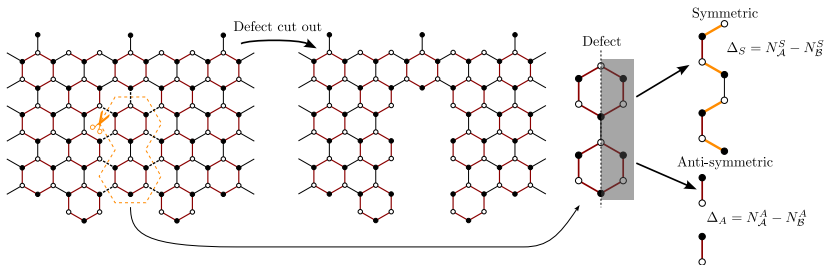
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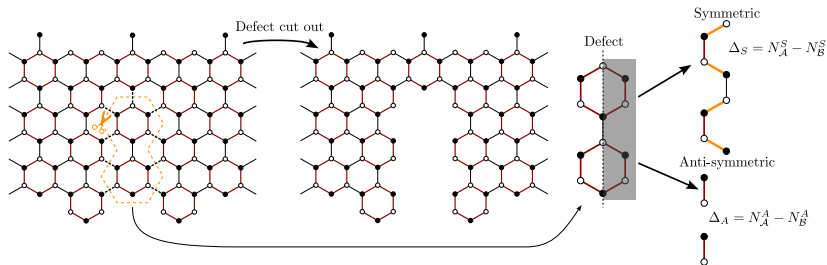
Perfect reflection

- When is it **perfect transmission**?
- Extra condition: **Defect must preserve topology**



Topological indices of defects

- For mirror symmetric and chiral defect
- Define **topological indices** (Δ_S, Δ_A)
- (Similar to mirror winding numbers)



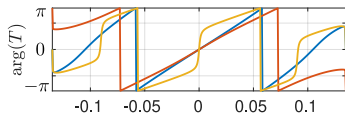
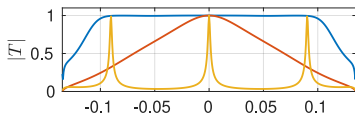
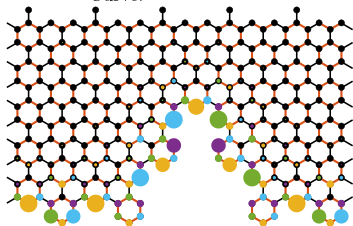
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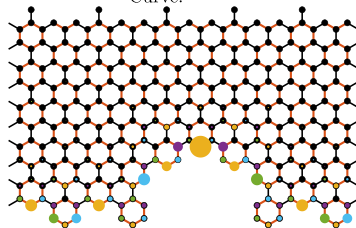
Nice result:

- If defect has $(0, 0)$ (trivial) always have $T(0) = 1$
- It is true for **any** defect with missing whole molecules

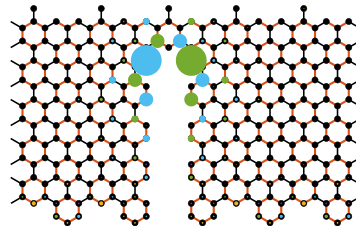
Transmission across trivial defects:

Curve: — ε 

Curve: —

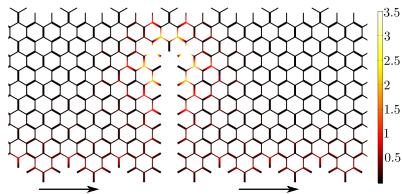
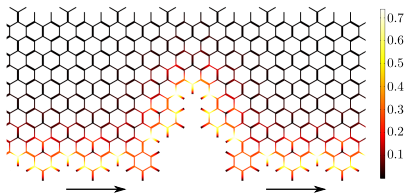
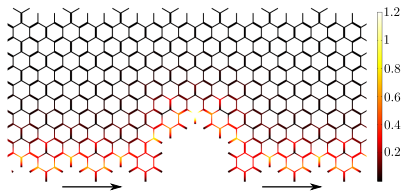
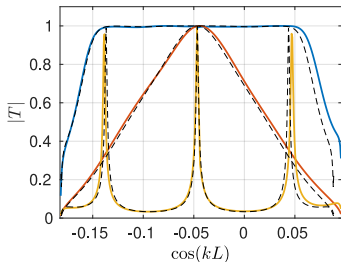


Curve: —



Acoustic realization: with networks of tubes

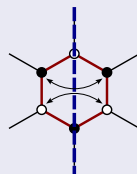
Dashed: lattice model
Color: 2D finite element simulations



Conclusion

Topological edge modes in Kekule:

- Combination of symmetries
- Chiral and mirror
- Edge dependent



Kekule ribbon with defects:

- Symmetries guarantee

$$R(0) = 0 \quad \text{or} \quad T(0) = 0$$

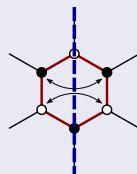
- Topological index: trivial defect $\Rightarrow T(0) = 1$

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Thank you.

Different types of edges:

